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# New Methods for Electrical Resistivity Measurement of Semi-conductive Powdered Materials<sup>\*</sup>

(Received October 22, 1954)

### Sanchi MIZUSHIMA\*\*

## Abstract

Two methods have been devised to measure the electric resistivity of powdered conductive materials. The one utilizes the eddy current effect in a conductor placed in an alternating magnetic field and the other is based on the principle that the conductivity of an electrolytic solution depends upon that of the powders of the specimen suspended in it. It is shown that both methods are practical enough.

## I. Introduction

In the course of investigation on the electric resistance of semi-conductors, especially of various kinds of carbons, it became necessary to devise some reliable methods which enable the measurement of the true resistance of powdered samples. Usually for such samples, only the apparent resistance is measured under determined pressure on a certain amount of the specimen: by extrapolating the resistance vs. pressure curve to infinite pressure the true resistance of the individual particles is determined. Since in this case the apparent resistance is composed of the bulk resistance and the contact resistance, between which the role of the latter exceeds that of the former always, this way of measurement is subject to some doubt, and moreover as there are many unknown or uncertain factors which seem to govern the measured values.

Raw materials for the production of carbon are frequently in the form of fine particles. Blocks of pitch or petroleum coke, the most important raw materials in the artificial graphite industry, are powdered into fine particles which are subsequently mixed with binding materials, and after being molded into desired forms are baked to the final products. The knowledge of the true resistance will be thus useful for the control of production as well as for the fundamental understanding of the changes of the electric and other physical properties during carbonization of organic materials and graphitization of amorphous carbons.

In the following sections are described two new methods for the measurement of the true electric resistance of powdered specimens. The one should not be said

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to be new in principle: it utilizes the eddy current loss phenomenon in individual particles in an alternating magnetic field. The other may be called substitution

particles in an alternating magnetic field. The other may be called substitution method: the sample is immersed in an electrolytic solution and the conductivity of the whole is measured, from the values of which the desired resistivity of the specimen is determined.

### II. Eddy Current Loss Method

The principle of this method is as follows. In a conductive material which is placed in an alterating magnetic field eddy current arises converting some amount of the applied field energy into heat. Since the lower the specific resistance of the sample the bigger the effect, the resistivity can be determined by the measurement of the losses. Now, when a spherical conductor of radius  $\alpha$  is in the field, the eddy current loss P is expressed by the equation, as long as the skin effect can be neglected,

 $P = (\pi/15)\omega^2(\alpha^5/\rho)Hm^2 \times 10^{-16}$  watt.,

where  $\omega$  is the circular frequency of the field,  $\rho$  is the specific resistance and Hm the peak value of the field. When the condition  $\alpha \sqrt{4\pi\omega/\rho} \ll 1$  is fulfilled, the skin effect is negligible.

Thus, the total loss of an assemblage of spherical particles is given by

 $P = (W/20D)\omega^2(\alpha^2/\rho)Hm^2 \times 10^{-16}$  watt.,

where D is the density of the bulk, not the apparent density, and W is the total mass.

If the magnetic field is produced by a solenoid and the specimen is placed in its central position, the effective resistance of the coil increases owing to the eddy current loss by the amount

 $R_{e} = (W/10D)\omega^{2}(\alpha^{2}/\rho)(0.4\pi N)^{2}/(l^{2}+4\pi r^{2}) \times 10^{-16}\Omega,$ 

where N,l and r are the number of windings, the length and the radius of the coil respectively. This relation makes the determination of the true resistance of individual particles possible provided that other quantities are known.

In the present experiment the increase  $R_e$  of the coil resistance was measured by a Q-meter. The quality Q of a coil is given by  $Q = \omega L/R$  at the series resonance with a variable condenser, where L is the inductance. R the resistance of the coil respectively.  $R_e$  is given then by

 $R_e = \omega L(Q'^{-1} - Q_0^{-1}),$ 

where  $Q_0$  and Q' are the values of Q for the cases without and with the specimen respectively.

Though the resonance frequency should vary more or less with the change of the coil inductance caused by the insertion of the specimen, the deresonance was found to be so small in this experiment that the additional adjustment was not needed throughout.

A number of coils were prepared, of which the construction data of a typical

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one is as follows. The full length is 11cm, the number of windings 8, pitch of the coil 1.4cm, radius 1.10cm. The bobbin is of a glass tube and the inductance  $0.588\mu H$ .

Because each particle has to be insulated from the other, the powdered samples were taken into a test-tube and filled with mineral oil. Moreover that the particles have sizes as equal as possible among them is desirable. They were separated into three stages by Tyler meshes No.  $30 \sim 45$ ,  $45 \sim 55$ , and  $55 \sim 80$ .

As for the shapes of these particles, it is natural that they are not perfectly spherical. But I contented myself in this experiment by the microscopic observation which revealed that they were rather block-like in form while needle-like or flake-like ones were lacking.

Now, preliminary tests were made on samples prepared by grinding and sieving from an artifical graphite electrode with the electric resistance of  $1.2 \times 10^{-3} \Omega \text{cm}$  in its original form. The following is the results obtained by the use of a  $0.472 \mu H$ coil.

Table	1.	Specific	resistance	of	artificial	graphite
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Particle size, mesh	Frequency MC	Specific resistance 10 <sup>-3</sup> Ωcm
No. 45~55	15 25	0.9
No. 55~80	15	0.9
//	25	1.1

Here the bulk density of the specimen was taken to be 2.0g/cm<sup>3</sup>.

As can be seen from the table, reasonable values are obtained, which are slightly less than the original value  $1.2 \times 10^{-3} \Omega$  cm.

Nextly the measurements were done on specimens of pitch coke. Blocks were baked for an hour at different temperatures given in Table 2 and then were powdered. The results shown there are for the experiments on particles ranging from mesh number 30 to 45 in size.

Baking temp. °C	Frequency MC	Specific resistance 10 <sup>-3</sup> 0cm
920	15	10
"	25	9
1100	15	6
"	15	6
industrial baking	15	4.5
//	20	5

Table 2. Specific resistance of pitch coke

In the calculation of the above table the bulk densities of these three specimens were taken to be 1.8, 1.9, and  $2.0g/cm^3$  respectively.

Experiments were tried also on samples with baking temperatures below 800°, which gave no results owing to the too high values of the resistivity.

Comparison of these tabulated values with the apparent ones of the original

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blocks was considered to be impractical since these blocks had numerous cracks and pores. The values in the table corresponding to different baking temperatures are, however, reasonable.

Table 3 shows the variation, if any, of the measured resistances with particle size. The specimens used therein are the same as the one given in the last columns in the above table.

s	Particle ize, iesh	Frequency MC	Specific resistance 10 <sup>-3</sup> Ωcm
No.	30~45	15	4.5
No.	45~55	"	5
No.	55~80	"	4
No.	80~110	//	4

Table 3. Measured resistivity vs. particle size

The speimen given in the last column of the above table is of particles with too small sizes that the loss is small and the result is less reliable.

It is not unexpected that the resistance varies with particle size because of the inherent porous nature of coke, but as is revealed by the table roughly equal values are obtained irrespective of particle size.

Thus, this method has been proved to be fairly practical for the measurement of the specific resistance of powdered carbons. If suitable selection of the frequency is made, the resistivity of many kinds of conductors will be able to be measured in this way.

# **III. Substitution Method**

Another method for the measurement of the specific resistance of fine particles was devised which is based on a different principle from that given in the previous section. In this method, the specimen in powdered form is immersed in an aqueous solution of electrolytes with known concentration and comparison is made between the apparent electric resistance of the solution and that of the pure so-

lution. Fig. 1 shows the measuring vessel, a part of the whole apparatus. As seen in the figure, the powdered specimen is dipped in a solution which fills the vessel and the conductance between the electrodes is measured by passing an alternating current. As long as the concentration of the solution and correspondingly its specific conductivity are high enough, the existence of the speci-



Fig. l.

men lowers the apparent conductance by a certain amount. When the concentration

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is subsequently decreased, since the conductivity of the solution itself decreases, there is one cross-point where the apparent conductivity coincides with that of the pure solution and the current passes as if there were no specimen in the solution.

Further lowering of the concentration results in the relative increase of the conductivity of the solution containing the specimen compared with that of the pure solution at the same concentration, as in this stage the conductivity of the specimen is higher than that of the solution. Thus, the measurement of the specific conductivity of the specimen is done by the determination of the point of coincidence of the conductance of the two solutions one of which contains the specimen while the other does not.

This method can be applied without recourse to the size, the form and the amount of particles of the specimen and moreover is not affected by the existence, if any, of insulating films on the particle surfaces. It should be noted, however, that the measurable range of specific resistance is limited because there are no suitable solutions which are applicable for specimens having specific resistance lower than about  $1\Omega$ cm.

For the purpose of studying the applicability of the above idea, tests were run on powders of several species of pitch coke baked at relatively low temperatures and of green carborundum. The particle size was so selected that they sink in the solution easily, and the specimen was poured into the vessel to occupy about  $60 \sim 80\%$  of the total current path by the apparent volume. As for the solution, aqueous solution of potassium chloride was used throughout while a conventional circuit composed of an A.C. bridge and a receiver was constructed to measure the conductance. The following is the results thus obtained.

Pitch coke 1; produced at the Wakamatsu factory of The Tokai Electrode Manufacturing Co. from coal pitch of the Kurosaki factory of The Mitsubishi Kasei Co. This was given laboratory-baking of  $620^{\circ} \times 30$ min. Results are given in Table 4.

Companyation N	Cond	uctance, µU
Concentration, N		Solution and Specimen
0.0025	29	35
0.005	58	61
0.01	120	110
0.02	240	160
0.1	1159	550

Table 4.	Pitch	colco	Chakad	<b>a</b> t	6200)	
Table 4.	FILCH	соке	Daked	ат	$n_2 (1^{-1})$	

Temperature 20°

Here the first row is for the concentration of the solution, the second for the conductance between electrodes in case of solution only and the third for that of solutions with specimens.

From the table the concentration at the cross-point is interpolated to be 0.007N which means that the specific resistance of the specimen is  $1.1K\Omega$ cm.

Pitch coke 2; this is obtained by baking at 700° for 45min. the same raw coke as the above. Results are given below.

Conductance, mu		
Pure solution	Solution and specimen	
0.09	1.75	
2.8	3.5	
8.3	8.3	
20.0	16.0	
	Pure solution 0.09 2.8 8.3	

Table 5. Pitch coke (baked at  $700^{\circ}$ )

From the table the concentration at the cross-point was found to be 1N and consequently the specific resistance of the specimen is  $10\Omega$ cm.

The above values of specific resistance of baked cokes are seen to be reasonable. Nextly similar measurements were done on the powders of "Erema" produced by The Tokai Konetsu Co., heating element composed mainly of silicon carbide, and also of the green carborundum which is used for the production of the former. Experimental results on Erema powders are given in Table 6.

Table V. TOwders of Diema	Table	6.	Powders	of	"Erema"
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Cor	nductance, $\mu \upsilon$
Pure solution	Solution and specimer
14	20
40	40
140	110
410	290
	Pure solution 14 40 140

From the table the cross-point was determined to be 0.003N and the specific resistance of the specimen was found to be  $2.5K\Omega$ cm.

For the powders of silicon carbide, a certain type of the measuring vessel was used in which the powder of the specimen was filled completely between the electrodes. In this case it was noticed that because of the high specific resistivity of the specimen the experiment had to be done in the small concentration range and consequently the conductivity of the distilled water which was used for the preparation of the solution came into play. But this does not result to any experimental errors. Results are shown in Table 7, where not the values of the conductance but those of the specific conductance are given.

Table 7.	Powders	of	silicon	carbide
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Concentration, N	Specific conductance, $\mu \upsilon cm^{-1}$		
Concentration, N	Pure solution	Solution and specimen	
0.00000	3.18		
0.00010	15.7	24.8	
0.00015	33.3	28.8	
0.0005	<b>6</b> 6.3	33.7	
0.0010	127.1	51.0	

Temperature 22°

In the table the specific conductance of the solution at zero concentration is just

that of the distilled water. From the position of coincidence the specific resistance of the specimen was determined to be  $35K\Omega$ cm.

Now, there seems to be no theoretical formula which expresses exactly the relations among the apparent specific conductivity  $\sigma_{app}$  of the solution containing the specimen, that of the pure solution  $\sigma_{sol}$  and that of the specimen itself  $\sigma_{sp}$ .

Making use of the formula for the apparent dielectric constant of a mixture of different dielectrics, we wish to approximate the actual dependency by

$$\sigma_{app} = \frac{2\sigma_{sol} + \sigma_{sp} - 2P(\sigma_{sol} - \sigma_{sp})}{2\sigma_{sol} + \sigma_{sp} + P(\sigma_{sol} - \sigma_{sp})} \times \sigma_{sol}$$

where P is the volume ratio of the powder of specimen against the whole volume.

Since this holds for the case of spherical particles and under the condition  $P \ll 1$ , the agreement between the theory and experiment is not to be expected to be good. When P is taken to be 0.71 the above equation agrees fairly well with



experimental results as seen in Fig. 2, while the measured value of P is 0.46. The curves in this figure corresponds to the data given in Table 6. But the agreement of this order of accuracy can be obtained simply by the equation

$$\sigma_{app} = (1 - P)\sigma_{sol} + P\sigma_{sp} .$$

In this case too, P has to be taken to be 0.46 to fit the data, showing the apparent conductivity of the whole is governed by that of the particles more strongly than expected from the volume ratio.

Thus the above processes enable the measurement of the specific resistance of powdered specimens in the range of  $1\sim10^{5}\Omega$ cm, where the following should be noted in the practical application.

Particles of the specimen should not react with the solution.
To cancel the capacitive reactance of the solution containing the specimen in the measurement, compensation by the use of condensers must be made in the corresponding arm in the A. C. bridge.

In conclusion, the author wishes to express his thanks to Mr. J. Okada of The Tokai Electrode Co. for kind collaborations.

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